

# STRUCTURAL SIMULATION OF CSSNBr3/CS2AGBIBR6-BASED PEROVSKITE HETEROJUNCTION SOLAR CELL

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Abstract— This article exhibits the performance analysis based on simulative investigation of a non-toxic, lead-free PIN type perovskite heterojunction solar cell (PTSC) with the inclusion of Cs<sub>2</sub>AgBiBr<sub>6</sub> as the intrinsic type absorber layer, CsSnI<sub>3</sub> or CsSnBr3 as the P-type absorber layers, CuSCN as the hole transport layer (HTL), TiO2 as electron transport layer (ETL), and CdS as Transparent conductive layer (TCO), respectively. From our first proposed structure, we attained an outstanding power conversion efficiency (PCE) of 37.98% when a CsSnI<sub>3</sub> as P-type absorber at optimum thickness of 300 nm was included. From the second proposed device structure, we achieved highest PCE of 37.77% when CsSnBr3 replaced the P-type absorber at the same optimum thickness of 300 nm. The device functioned optimally when the electron affinities were set at 3.2 eV, 4.190 eV, 4.470 eV and 3.4 eV for n-TiO<sub>2</sub>, Cs2AgBiBr6, CsSnI3 and CuSCN, respectively. We observed the effect of thickness and doping density on the device's performance of different charge transport layers. Inclusively, this simulation results will render perceptive guidance towards producing an environmentally benign PTSC through pertinent observation of the transport mechanism and understanding of the role of a suitable TCO and an HTM on the routine features of PIN type perovskite heterojunction solar cell.

*Keywords*— Cs<sub>2</sub>AgBiBr<sub>6</sub>, CsSnI<sub>3</sub>, CsSnBr<sub>3</sub>, perovskite heterojunction solar cell (PTSC), ETL, HTL, PCE, TCO.

#### I. INTRODUCTION

The world is now grappling with the burning need of clean and renewable energy driven power production only for the demand of the elimination of greenhouse gas emission generated from fossil fuel based energy consumption. Among the accessible renewable-energy sources, solar power is considered as the most plenteous particularly for South Asian countries due to adequate sun irradiance. Hence, we can see that the use of photovoltaic (PV) solar cells is on the progress and certainly, their application will play significant role in future sustainable energy systems.

Although currently in terms of commercial aspect Sibased PVs are the most commonly used cells, however they render a few major limitations such as restricted efficiency and demanding cost [1-6]. Therefore, some other evolving PV technologies [7] such as organic PV, hybrid organic-inorganic PV cells [8,9], dye-sensitized PV cells [10], quantum dots or low-dimensional heterostructure based PVs [11] and perovskite PV cells [12,13] which have appealed the scientists owing to their promising features. Such as, Halide based perovskites (ABX<sub>3</sub>) are gaining attention because of suitable light capturing attribute and considered as third generation of PV. They have also created applications as memory devices, LED's, LASER's and detectors [14-16]. In overall, this perovskites exhibit cubic structure with ABX<sub>3</sub> composition, where A = organic monovalent cation (for example CH3NH3+, CH3(NH2)2+), B = metal divalentcation (for instance Pb2+, Sn2+, or Ge2+) and X may belong to Halogen family (X = I, Br, Cl). Recently, organic-inorganic halide perovskites have attained substantial interest as potential candidates in the field of photovoltaics industry due to its favorable optoelectronic properties such as enhanced absorption-coefficient, longer electron and hole diffusion length, suitable charge carrier mobility, recombination lifetime, and lesschallenging fabrication procedures [8,9,12,13]. Within less than a decade, the PCE of perovskite PV cells have reached a record more than 25%, exceeding the performance of conventional CdTe/Cu(In, Ga)Se<sub>2</sub>-based

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thin-film PV cells [19], whereas, the tandem metal halide perovskites have achieved more than 26% of efficiency, surpassing the crystalline silicon [20–21]. Nevertheless, the organic elements (MA and FA) causes poorer instability under several ambient conditions and thus compromises with the overall performance [22-24]. It is often proven that replacing these organic counterparts by inorganic alkali cations such as Cs+ or Rb+ can resolve this problem [25,26]. Recently, inorganic CsPbX<sub>3</sub> perovskites are investigated as a budding candidate for efficient PSCs due to better thermal stability. However, although the Pb-based perovskites offers high figure of merits, their viable application is mired due to the moist and heat instability, and together with the toxicity of lead-which cause serious damage in the environment and can cause low and premature birth rates, cancer in circulatory system, brain-cell damage and kidney failure. These urges the researchers to give emphasis on Pb-free, non-toxic alternatives. Recently, substitution of Pb with Sn and Ge is found reliable due to comparable electronic features [27]. Sn seems fairly suitable; particularly acts as favorable candidate for bulk scale production owing to the longer oxidation state, appropriate absorption coefficient (>105 cm<sup>-1</sup>), finest forbidden-gap (1.2-1.62 eV) and well air stability [28-30]. So far, the maximum efficiency of lead-free perovskite based solar cells is achieved 13.24% for tine-based structure [31]. It was reported that perovskite materials (ABX<sub>3</sub>) with Br in the X-site position could render greater stability compared to its I equivalent [32]. With the exchange of I with Br, orthorhombic phase of CsSnI<sub>3</sub> converts into cubic phase of CsSnBr<sub>3</sub>, and can render better electronic properties such as increased open circuit voltage in mixed halides [32]. In line with this, halide double perovskite expressed by the form of A2M'M"X6 (A is a monovalent cation, M' and M" are monovalent and trivalent metal ions), has also become favorable alternatives. It also offers high environmental stability, suitable optoelectronic properties, reduced toxicity, and can be considered as good substitute for Pbbased perovskites. For example, lead-free, Bi-based Cs<sub>2</sub>AgBiBr<sub>6</sub> halide double perovskite can be a good candidate for PV solar cell. Its thermal and atmospheric stability, lengthy carrier lifetime, better absorption coefficient, and suitable bandgap (ranging from 1.83 eV to 2.19 eV) make it proper as the absorber layer [33]. Too, this perovskite crystallizes in the elpasolite structure, absorbs light in the visible range of the solar radiation [33]. Even though these all-inorganic perovskite based solar cells have shown remarkable growth in the recent years, however, there remains substantial concern in increasing the merits of all-inorganic perovskite solar cells through collective optimization of the absorber layers and its interface.

In this work, we have systematically studied the structural, electronic, and optical properties of two slightly different heterojunction structures with lattice matched HTL and ETL having apposite electron affinities. Besides, the photovoltaic performance of our halide perovskites ( $Cs_2AgBiBr_6/CsSnI3$  and  $Cs_2AgBiBr_6/CsSnBr_3$ ) has been observed by varying the thickness and doping concentrations.

## II. STRUCTURAL SPECIFICS AND SIMULATION METHOD OF PROJECTED PHOTOVOLTAIC CELLS

In this simulation, we proposed two slightly different devices having the heterojunction architecure of n-type CdS/n-typeTiO2/I-type Cs2AgBiBr6/p-type CsSnI3/ptype CuSCN and n-type CdS/ n-type TiO2/I-type p-type Cs<sub>2</sub>AgBiBr<sub>6</sub>/ CsSnBr3/p-type CuSCN. respectively. We witnessed the effect of absorber thickness, the doping density, absorber defect density, radiative recombination rate (band-to-band), the electron affinity of HTL and ETL, back contact work function, on the overall performance of our device. We are hopeful that these simulation results will contribute to the laboratory based fabrication of non-toxic, and viable halide based PTSC.

## A. Selection of HTL, ETL and absorbers for the devices

The inclusion of electron/hole-transport materials (ETL/HTL) in the PTSC is very crucial for extracting photo-generated excitons from the perovskite absorber layer to the outer contacts. Selection of ETL and HTL is vital for active carrier transport at the interface mostly due to the energy band arrangement; the ETL pulls electrons while obstructing the holes and HTL pulls holes by obstructing holes. The mostly met ETL materials in PVSCs are Fe<sub>2</sub>O<sub>3</sub>, SnO<sub>2</sub>, ZnO, TiO<sub>2</sub>, In<sub>2</sub>O<sub>3</sub>, a-IGZO and PCBM [34–36]. Among these, TiO<sub>2</sub> is found as the most usually and oftenly used ETL owing to its some attractive attributes such as low-cost, high electron mobility, low-ecotoxicity and stability.

The utmost common material for HTL are Spiro-OMeTAD, PEDOT:PSS, P3HT, and PTAA. However, due to their costly fabrication procedure and low chemical durableness many researches have been trying to find some other auxiliary. In addition, some inorganic HTL (Cu-based) such as CuSCN, CuI, Cu<sub>2</sub>O, and CuOx are proposed as relevant materials owing to their quicker hole mobility, chemical durability and moderate-fabrication-cost [37,38].

Likewise, insertion of effective absorber layers in tandem cells and their proper alignment of energy bands with transport layers [10] is one of the prime concerns for effective passage or collection of the charge carriers to the metallic-contacts. This is basically possible if the work-



function of front contact material is nearly matching or marginally larger than conduction band edge of the transport layer. Consequently, higher work-function presumably yields an ohmic contact for holes and electrostatic barrier for electrons in HTL. Cell performance also deteriorates if the recombination mechanism is not addressed properly as gathered charges will influence the electric field inside the cell. Thus appropriate buffer layer is vital to maintain the current flow through an entire tandem architecture. In light of the above issues, we have identified Cs<sub>2</sub>AgBiBr<sub>6</sub> as an thin intrinsic (I) layer and act as low dimensional quantum confined absorber layer that can capture more excitons in the vicinity of wide p-type absorbers (CsSnI<sub>3</sub> or CsSnBr<sub>3</sub>). This Cs<sub>2</sub>AgBiBr<sub>6</sub> is all-inorganic, lead-free perovskite, which adopt the ideal cubic structure with very close lattice parameters and particularly lattice matched with CgSnBr<sub>3</sub> at room temperature. Such combination minimizes interference defect states from complex compositions of crystal structures.

## B. Simulation method for the devices

well-known SCAPS-1D The simulation program/software (version 3307) is used in this study to carry out the modeling and simulation. SCAPS-1D is a one-dimensional simulator facilitating maximum seven semiconductor layers and it tasks based on Poisson and Continuity equations to model and simulate numerous PV structures. This simulator can calculates energy bands, concentrations and currents, J-V characteristics, spectral response (QE). This solar cell simulator has been widely used for the device modeling of solar cells and results almost validate the experimental results. The calculation was done under 1.5 AM solar irradiance and at a temperature of 300 K.

The all requisite design parameters used for the device modelling of the structures are given in Table-I and II; collected from formerly reported literatures [33,39].

In Figure 1 (a), the structural details of our proposed PIN type two different structures are shown pictorially. In our first proposed structure, CdS is inserted as n-type (optimized 10 nm thick) at the top,  $TiO_2$  is used as n-type (optimized 20 nm thick) ETL,  $Cs_2AgBiBr_6$  is utilized as an intrinsic type absorber (optimized 10 nm) and the p-type CsSnI<sub>3</sub> (300 nm thick) is sandwiched as the absorber layer. Thereafter a p-type CuSCN (optimized at 20 nm) is inserted as HTL. The total thickness of the entire structure becomes only 360 nm as a whole. From the point of view of fabrication suitability due to favorable lattice matching, we studies another structure that is almost similar; only the p-type CsSnI<sub>3</sub> is replaced by the p-type CsSnBr<sub>3</sub> shown in Figure 1 (b).



(b)

Fig. 1. (a) Schematic depiction of the first proposed structure [lead-free PIN type perovskite heterojunction solar cell (PTSC) with the inclusion of  $Cs_2AgBiBr_6$  as the intrinsic type absorber layer,  $CsSnI_3$  as the P-type absorber layers, CuSCN as the hole transport layer (HTL), TiO<sub>2</sub> as electron transport layer (ETL), and CdS as Transparent conductive layer (TCO)].

(b) Schematic representation of the second PTSC with the inclusion of  $Cs_2AgBiBr_6$  as the intrinsic type absorber layer, CsSnBr3 as the P-type absorber layers, CuSCN as the hole transport layer (HTL), TiO<sub>2</sub> as electron transport layer (ETL), and CdS as Transparent conductive layer (TCO)

PARAMETERS OF KEY MATERIALS UTILIZED AS ABSORBER LAYER						
Parameter	$Cs_2AgBiBr_6$	CsSnI <sub>3</sub>	CsSnBr <sub>3</sub>			
Thickness (nm)	10 nm	300 nm	300 nm			
$E_g(eV)$	2.05	1.270	1.110			
$\chi \left( eV\right)$	4.190	4.470	4.460			
$\epsilon_r$	5.800	10.590	6.400			
$N_c(cm^{-3})$	$1.00 \times 10^{20}$	$1.8 \times 10^{19}$	$1.5 \times 10^{19}$			

TABLE I ARAMETERS OF KEY MATERIALS UTILIZED AS ABSORBER LAYER



$N_v(cm^{-3})$	$1.800 \times 10^{19}$	$1.67 \times 10^{18}$	$1.47 \times 10^{18}$
V <sub>e-th</sub> (cm /s)	$1 \times 10^{7}$	$1 \times 10^7$	$1 \times 10^{7}$
$V_{p-th}(cm/s)$	$1 \times 10^{7}$	$1 \times 10^{7}$	$1 \times 10^{7}$
$\mu_c(cm^2/V_{ m S})$	$1.181 \times 10^{1}$	$3.11 \times 10^{2}$	$4.37 \times 10^{2}$
$\mu_p(cm^2/V_{ m S})$	$4.90 \times 10^{-1}$	$4.2 \times 10^{1}$	$4.37 \times 10^{1}$
$N_D(cm^{-3})_{.}$	$1 \times 10^{19}$	0	0
$N_A(cm^{-3})_{.}$	1 × 10 <sup>19</sup>	5 × 10 <sup>19</sup>	$5 \times 10^{19}$

TABLE II					
PARAMETERS OF MATERIALS USED AS ETM, TCO AND HTM					

Parameter	n-TiO <sub>2</sub>	n- CdS	p-CuSCN
Thickness (nm)	20 nm	10 nm	20 nm
$E_g(eV)$	3.2	2.4	3.4
$\chi \left( eV\right)$	4.2	4.2	2.1
$\epsilon_r$	9	9	10
$N_c(cm^{-3})$	$2.2 \times 10^{18}$	$2.2 \times 10^{19}$	$2.5  imes 10^{18}$
$N_v(cm^{-3})$	$1.8 \times 10^{19}$	$1.8 \times 10^{18}$	$1.8  imes 10^{19}$
V <sub>e-th</sub> (cm /s)	$1 \times 10^7$	$1 \times 10^{7}$	$1 \times 10^7$
V <sub>p-th</sub> (cm /s)	$1 \times 10^7$	$1 \times 10^{7}$	$1 \times 10^{7}$
$\mu_c(cm^2/V_{ m S})$	$2 \times 10^{1}$	$1 \times 10^{2}$	$2 \times 10^{-4}$
$\mu_p(cm^2/V_{ m S})$	$1 \times 10^{1}$	$2.5 \times 10^{1}$	$1 \times 10^{-2}$
$N_D(cm^{-3})_{.}$	$1 \times 10^{18}$	$3.0 \times 10^{18}$	0
$N_A(cm^{-3})$	0	$1 \times 10^{1}$	$1 \times 10^{18}$

### III. RESULTS AND DISCUSSIONS

An effective heterojunction cell demands a few conditions, such as: the short circuit throughout the multijunctions should be properly matched and the PV will offer smallest, generation of excitons against wide range of solar irradiance should be effective while the electrical resistance and optical losses should be as low possible. These necessities can be controlled by varying band-gap, thickness, and by including suitable transport layer with suitable affinity levels.

It has been known that the PV cells typically gives higher efficiency if light enters into the cell from P-type window. As our structure is a PIN type we have chosen ptype-CsSnX<sub>3</sub> (X= I or Br) as the prime absorber layer which absorbs mostly the photons of visible to near infrared wavelength. The intrinsic (I) layer works as additional absorber layer that can capture multiple photons of higher energies due to its confined energy states attributed to 10 nm thickness and will also offer high electric field due to small thickness. On the other hand, wider p-type-CsSnX3 will soak more carriers due its wider thickness. Then the generated carriers experience the field, swept towards the outer contact of the structure pulled by the ETL/HTL. As a result, it is projected that the thickness, energy gap, and defect states of the I-layer are crucial design characteristics.



Fig. 1. The PCE as a function of varying thickness from 10 nm to 90 nm for (a) TCO (CdS) and (b) ETL (TiO2).





Fig. 2. The PCE as a function of thickness from 10 nm to 90 nm for  $Cs_2AgBiBr_6$  as I type absorber layer.

Our simulation, as shown in Figure 1, shows the effect of varying thickness on the performance of the cell; exemplarily the power conversion efficiency (PCE) is depicted. We can see that the optimized thickness of TiO<sub>2</sub> is achieved as 20 nm and for CdS the suitable thickness is found as 10 nm. The maximum PCE is achieved 39.88% with FF of 83.12%, open circuit voltage V<sub>oc</sub> of 1.17 V and short circuit current  $J_{sc}$  of 39.19 mA/cm<sup>2</sup>. Table III and IV depicts the obtained simulation result for ETL (TiO2) and TCO (CdS).

Table III							
Performance optimization of PTSC for ETL, TiO <sub>2</sub>							
	Voc J <sub>sc</sub>						
Thickness	(Voltage)	(mA/cm <sup>2</sup> )	FF (%)	PCE%			
0.01	1.17	38.27	83.11	37.16			
0.02	1.17	39.19	83.12	37.98			
0.03	1.17	37.25	81.98	36.59			
0.04	1.17	37.2	81.99	35.55			
0.05	1.17	37.16	81.99	35.52			
0.06	1.17	37.13	82	35.48			
0.07	1.17	37.09	82.01	35.48			
0.08	1.165	37.06	82.01	35.42			
0.09	1.165	37.02	82.02	35.39			

Table IV

Performance optimization of PTSC for TCO, CdS					
	J <sub>sc</sub>				
Thickness	Voc (Voltage)	(mA/cm <sup>2</sup> )	FF (%)	PCE%	
0.01	1.17	39.19	83.12	37.98	
0.02	1.17	38.27	82.22	36.63	
0.03	1.17	36.47	81.99	34.84	
0.04	1.16	35.77	81.97	34.15	
0.05	1.16	35.15	81.95	33.54	
0.06	1.16	34.61	81.93	33	
0.07	1.16	34.12	81.91	32.52	
0.08	1.16	33.69	81.9	32.09	
0.09	1.16	33.31	81.88	31.72	

Similarly, we can see from the Figure 2 that the maximum PCE is obtained at 10 nm thickness for  $Cs_2AgBiBr_6$  absorber layer (Table V shows the details).

Table V

Performance optimization of PTSC for Cs <sub>2</sub> AgBiBr <sub>6</sub>					
	Voc		FF		
Thickness	(Voltage)	$J_{\rm sc}~({\rm mA/cm^2})$	(%)	PCE%	
0.01	1.17	38.27	83.11	37.98	
0.02	1.17	37.27	82.02	35.63	
0.03	1.17	36.55	80.97	34.51	
0.04	1.17	36.02	79.99	33.6	
0.05	1.17	35.6	79.05	32.84	
0.06	1.17	35.26	78.16	32.17	
0.07	1.17	34.96	77.32	31.57	
0.08	1.17	34.71	76.53	31.04	
0.09	1.17	34.48	75.79	30.55	

On the contrary, we see from the Figure 3 that the PCE remains almost constant until 70 nm of thickness for HTL (CuSCN). We chose a thickness of 20 nm from the viewpoint of cost optimization.



Fig. 3. The effect of the variation of thickness on (a) PCE and (b) FF of the entire structure for the HTL (CuSCN).







Fig. 4. The variation of doping concentration on the performance parameters such as: (a) Fill-Factor (FF), (b)  $J_{sc}$  (short circuit current density) (c) open circuit voltage  $V_{oc}$ , and (d) efficiency (%) of the cell by varying the thickness of the p-type CsSnI<sub>3</sub> as the primary absorber layer.

The efficiency grows linearly when the thickness of the absorber I-layer is increased starting at 100 nm and eventually reaches its maximum at 300 nm, according to simulation findings. Next, we examine the impact of thickness, defect, and doping densities on our structure (Figure 4). After then, the efficiency starts to gradually decline. Yet, it is also recognised that much bigger thickness will disturb the electric field, therefore we find that the 300 nm p-absorber layer of the PIN structure is the optimal thickness. As we know, the photon absorption is projected to rise with increasing thickness.

Lastly, we also explore the similar effect of thickness of all layers in the second proposed structure and the performance parameters of only the main absorber layer  $CsSnBr_3$  is depicted in Figure 5. We can see from the analysis that the overall PCE is compromised by only 0.2% with an prominent increase of FF by 3% inferring reduction of series resistance in the entire structure. These findings suggest that  $CsSnBr_3$  can be a preferable alternative of  $CsSnI_3$  because it can more fabrication friendly for our proposed I layer of  $Cs_2AgBiBr_6$ considering lattice matching.





(%) ) of the cell by varying the thickness of the p-type  $CsSnBr_3$  as the absorber layer.

Lastly, the screenshot of the I-V curve from SCAPS-1D following the simulation of the second proposed PTSC, the calculated energy diagram, the calculated recombination current densities and the internal quantum efficiency curve are shown in Figure 6, Figure 7 and the plotted I-V curve of the same structure is shown in Figure 8, respectively for the readers understanding.

A Table VI compares the key performance merits of  $Cs_2AgBiBr_6/CsSnI_3$  and  $Cs_2AgBiBr_6/CsSnBr_3$  based device acquired in this work, as well as with those reported previously. We can see that our structure offers outstanding performance with the inclusion of lead free absorbers in both PTSCs. We hope this will motivate the researchers who are interested in perovskite based new generation of solar cell design.

COMPARISON WITH PREVIOUS WORK				
Device	Voc	J <sub>sc</sub>	FF (%)	PCE
	(V)	(mA/c		(%)
		$m^2$ )		· · /
		,		
CdS/TiO <sub>2</sub> /Cs <sub>2</sub> AgBiBr <sub>6</sub>	1.17	39.19	83.12	37.98
/CsSnI <sub>3</sub> /CuSCN (our				
proposed device 1)				
CdS/TiO2/Cs2AgBiBr6	1.03	37.78	85.94	37.77
/CsSnBr <sub>3</sub> /CuSCN (our				
proposed device 2)				
ZnO-NR/Cs <sub>2</sub> AgBiBr <sub>6</sub> /	0.92	11.10	44.02	4.48
P3HT (Base) [33]	0.72			
19111 (Buse) [55]				
FTO/ZnO/Cs2AgBiBr <sub>6</sub> /Ni	1.29	20.69	81.72	21.88
O/Au [40]	V			

 TABLE VI

 COMPARISON WITH PREVIOUS WORK



Fig. 5. The variation of doping concentration on the performance parameters (a) Fill-Factor (FF), (b)  $J_{sc}$  (short circuit current density) (c) open circuit voltage  $V_{oc}$ , (d) and efficiency

Fig. 6. The screen-shot image of the I-V curve from SCAPS-1D after simulation of the second proposed PTSC







Fig. 7. (a) The calculated energy diagram. (b) The calculated recombination current densities and (c) The quantum efficiency curve.



Fig. 8. The I-V characteristics of the PIN type PTSC for  $Cs_2AgBiBr_6/CsSnBr_3$ .

#### IV. CONCLUSION

In this paper, a methodical analysis two types of tandem PIN devices n-CdS/n-TiO2/I-Cs2AgBiBr6/p-CsSnI3/p-CuSCN and n-CdS/ n-TiO2/I-Cs2AgBiBr6/p-CsSnBr3/p-CuSCN, PV cells are modelled and calculated using SCAPS-1D simulator. The devices were simulated and optimized to gain best figure of merits by varying thickness, defect densities and doping densities of all layers. The highest PCE is achieved as 37.88% implementing Cs<sub>2</sub>AgBiBr<sub>6</sub> as thin I-type absorber layer with CsSbI<sub>3</sub> as suitable wide P-type exciton generation layer. The another structure results in almost same PCE of 37.77% by replacing the CsSbI<sub>3</sub> with CsSbBr<sub>3</sub>. Since Cs<sub>2</sub>AgBiBr<sub>6</sub> is an ideal, cubic structure with very close lattice matched with CsSnBr3 at room temperature. This combination of Cs<sub>2</sub>AgBiBr<sub>6</sub> with CsSnBr<sub>3</sub> certainly minimizes interference defect states from complex compositions of crystal structures. The findings of the simulation show that they played a key role in reaching such a high efficiency.

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